

# 5

# Charge Formation and Delocalization

# Table of Contents

## 1. Introduction

- 1.1. Brief History of Organic Electronics
- 1.2. Challenges in Organic Electronics

## 2. Intramolecular Electron Delocalization

- 2.1. A Primer on Quantum Mechanics
- 2.2. From Atomic Orbitals to Molecular Orbitals
- 2.3. Molecular Orbitals in  $\pi$ -Conjugated Systems

## 3. Electron Delocalization in Organic Materials

- 3.1. The Origin of  $\pi$ - $\pi$  Interactions
- 3.2. Organic crystals of  $\pi$ -Conjugated Molecules
- 3.3. Intermolecular Electron Delocalization

## 4. Intrinsic and Extrinsic Electronic Perturbations

- 4.1. Vibronic Coupling
- 4.2. Light-Matter Interaction

## 5. Charge Formation and Delocalization

- 5.1. Solitons
- 5.2. Polaron
- 5.3. Charges in Organic Materials
- 5.4. Charges at Interfaces

## 6. Charge Transport in Organic Materials

- 6.1. Overview of Transport Regimes
- 6.2. Band and Band-Like Transport
- 6.3. Polaronic Transport
- 6.4. Disorder-Controlled Transport
- 6.5. Towards a Unified View

## 7. Basic Organic Electronic Devices

- 7.1. Organic Field-Effect Transistors
- 7.2. Organic Photovoltaic Devices
- 7.3. Organic Light-Emitting Diodes

## 8. Organic Semiconductor Materials Preparation

- 8.1. Synthesis of  $\pi$ -Conjugated Molecules
- 8.2. Preparation of Thin Films
- 8.3. Patterning for Devices

## 9. Advanced Topics

- 9.1. Organic Sensors
- 9.2. Bioelectronics
- 9.3. Sustainable Electronics
- 9.4. Magneto-optoelectronic Devices
- 9.5. Spintronics

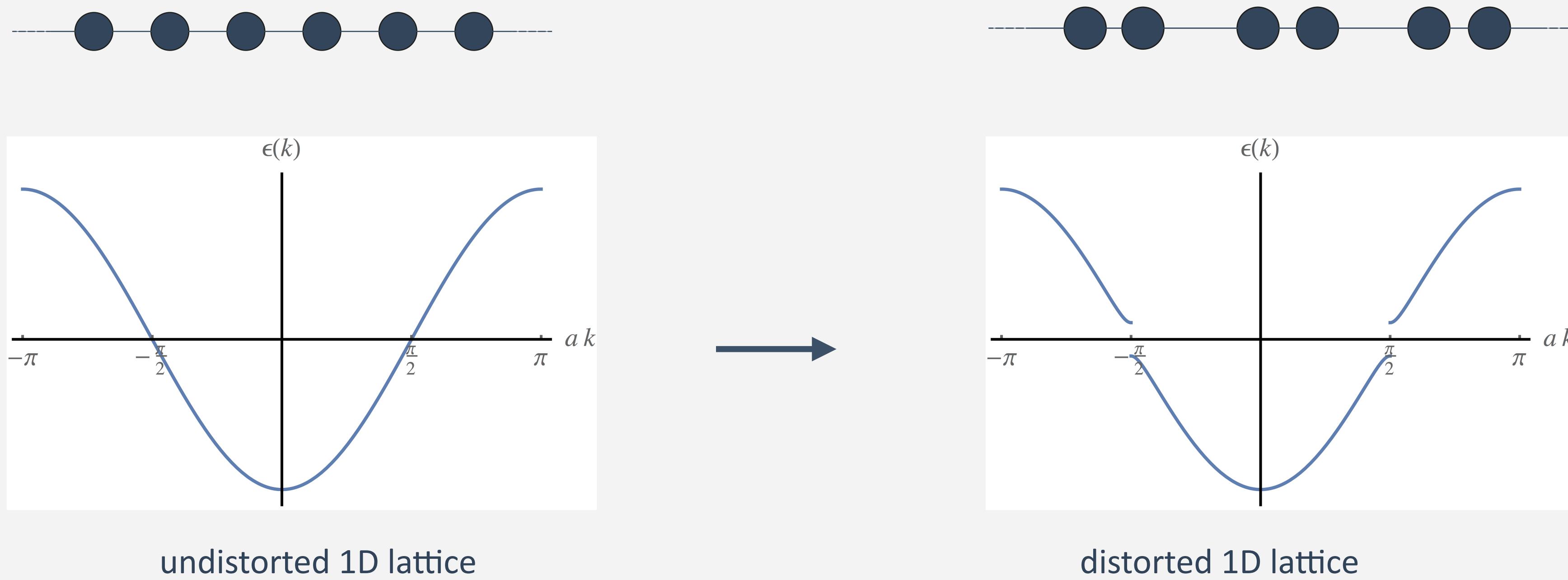
---

## 5.1 Solitons

---

# Polyacetylene Degeneracy: Peierls Instability

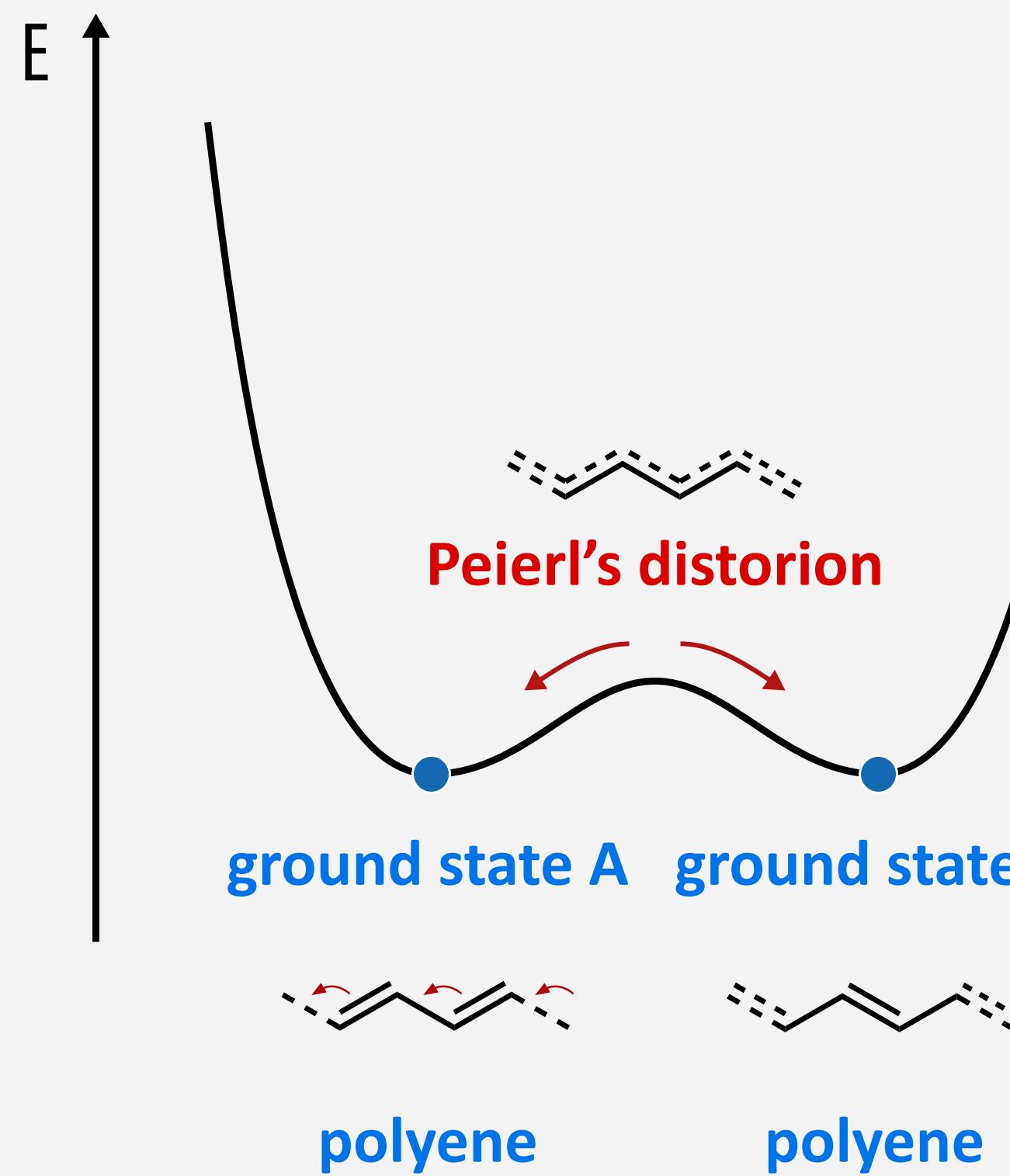
- Peierls' theory: one dimensional equally spaced lattice with one electron per ion is unstable



- solving the Schrödinger equation for a periodic potential in one dimension results in energetically more stable solution for bond length alternation
- energy gained by opening a band gap and thus lowering the filled energy levels is larger than the energy lost by the dimerization of carbon atoms

# Degenerate Ground States In Poly(acetylene)

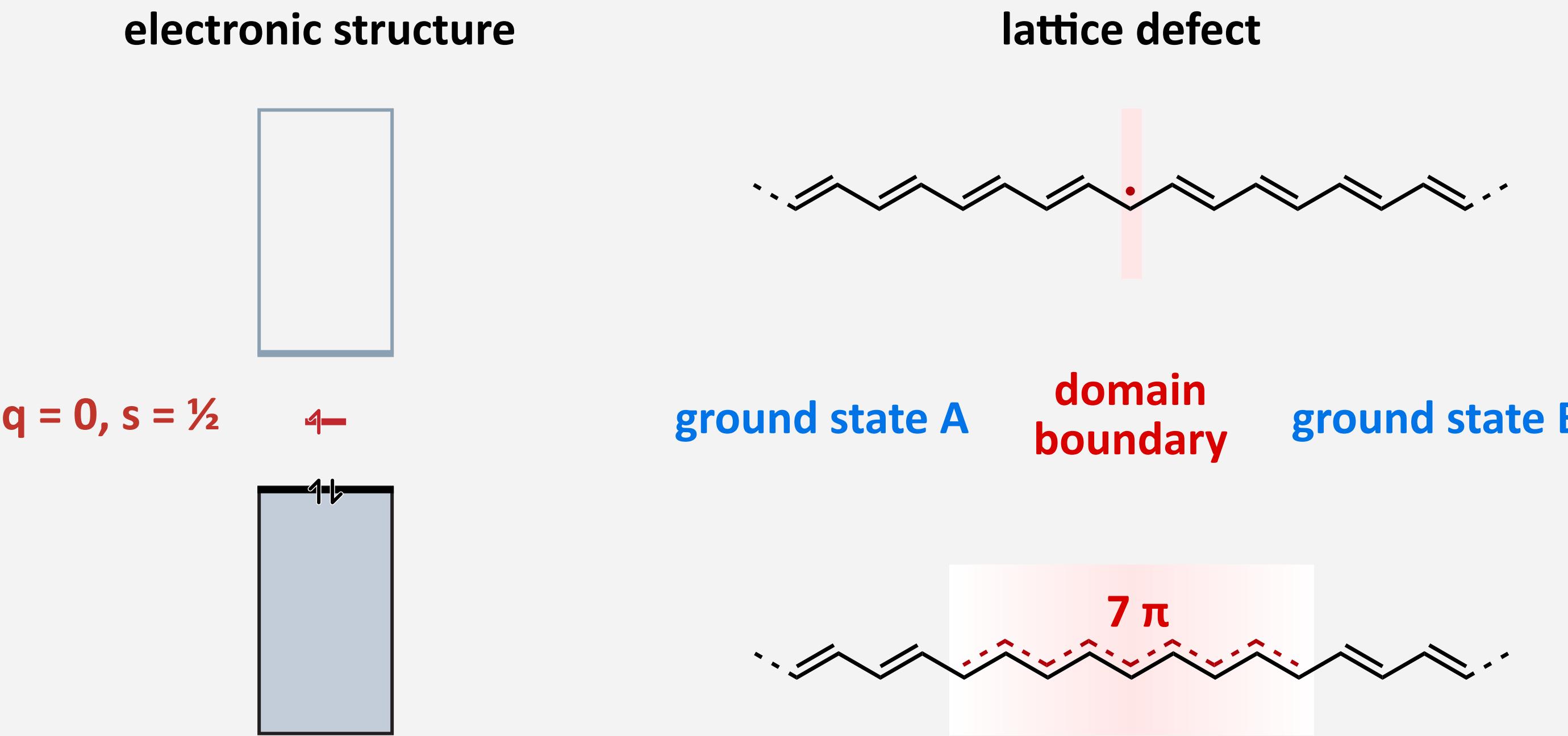
- poly(acetylene) is the only  $\pi$ -conjugated polymer with two equivalent resonance structures



- poly(acetylene) has **two energetically & geometrically equivalent, degenerate ground states**
- poly(acetylene) is unique in the sense that it has solitons as charge carriers

# “Spontaneous” Formation of Neutral Solitons in Poly(acetylene)

- at finite temperature, “lattice defects” at domain boundaries result in **neutral solitons**



- lattice distortion results in defect energy level in band gap with limited delocalization**
- neutral solitons are particles with spin  $s = \frac{1}{2}$  and charge  $q = 0$  plus lattice distortion**
- organic chemistry view: neutral solitons are radicals**
- neutral solitons cannot contribute to conduction, but are easily oxidized/reduced**

# SSH Model of the Soliton

- Electron Spin Resonance (ESR) spectroscopy (based on Zeeman splitting  $\Delta E = m_s g_e \mu_B B_0$ )
  - $g$  factor of  $g_e = 2.0026$ , similar to a free electron (2.0023)
  - spin density of  $1.23 \cdot 10^{19} \text{ cm}^{-3}$  or 1 soliton per 3200 carbon atoms (vs  $\text{MnSO}_4 \cdot \text{H}_2\text{O}$ )
  - spin density insensitive to  $\text{NH}_3$  vapor (used to quench inadvertent charge carriers)
  - Lorentzian peak shape, no hyperfine structure, line narrowing at higher temperatures imply delocalization and thermally activated mobility of the soliton defect

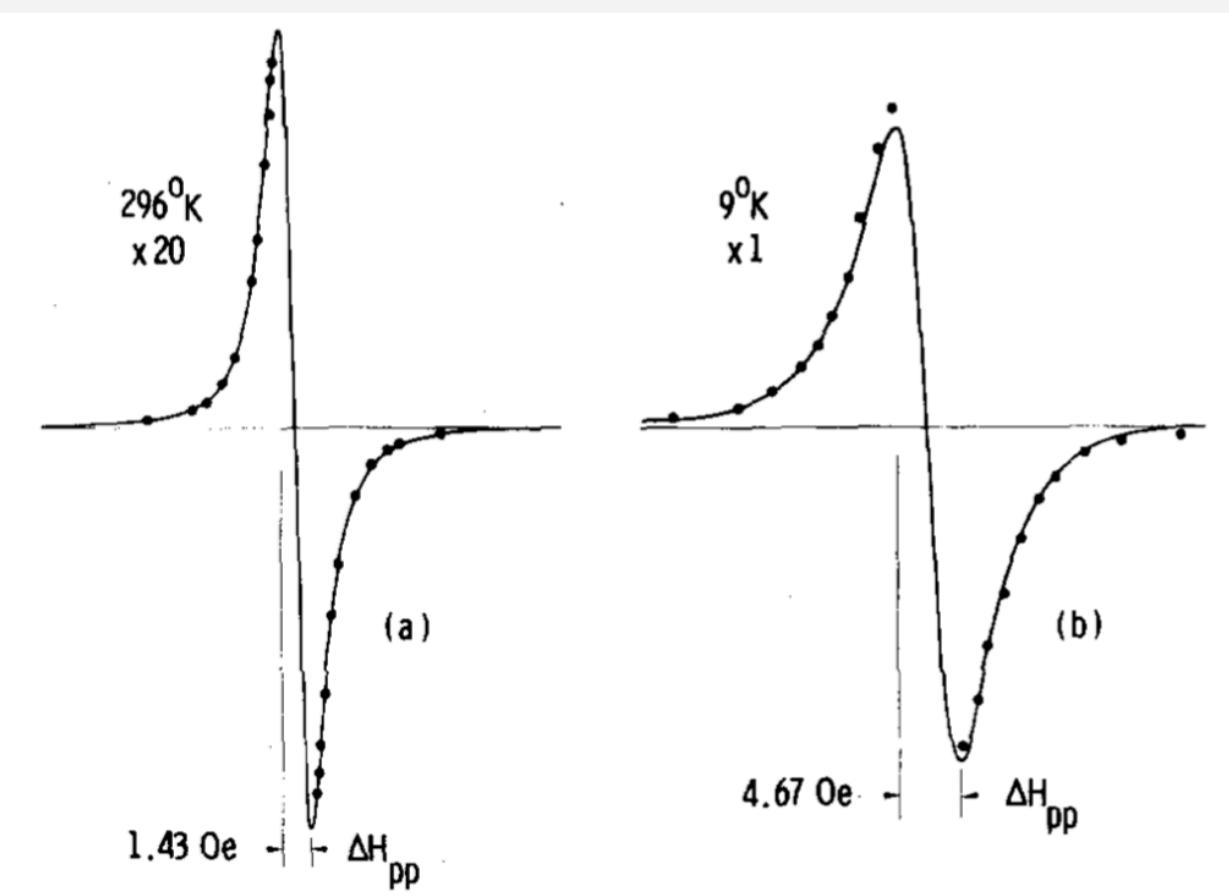


FIG. 1. ESR spectra (first derivative) of films of *trans*-(CH)<sub>x</sub> at 296 °K (a) and 9 °K (b). The circles show the theoretical lineshape based upon the peak-to-peak linewidth of the ESR spectra. (Magnetic field scales are not identical.)

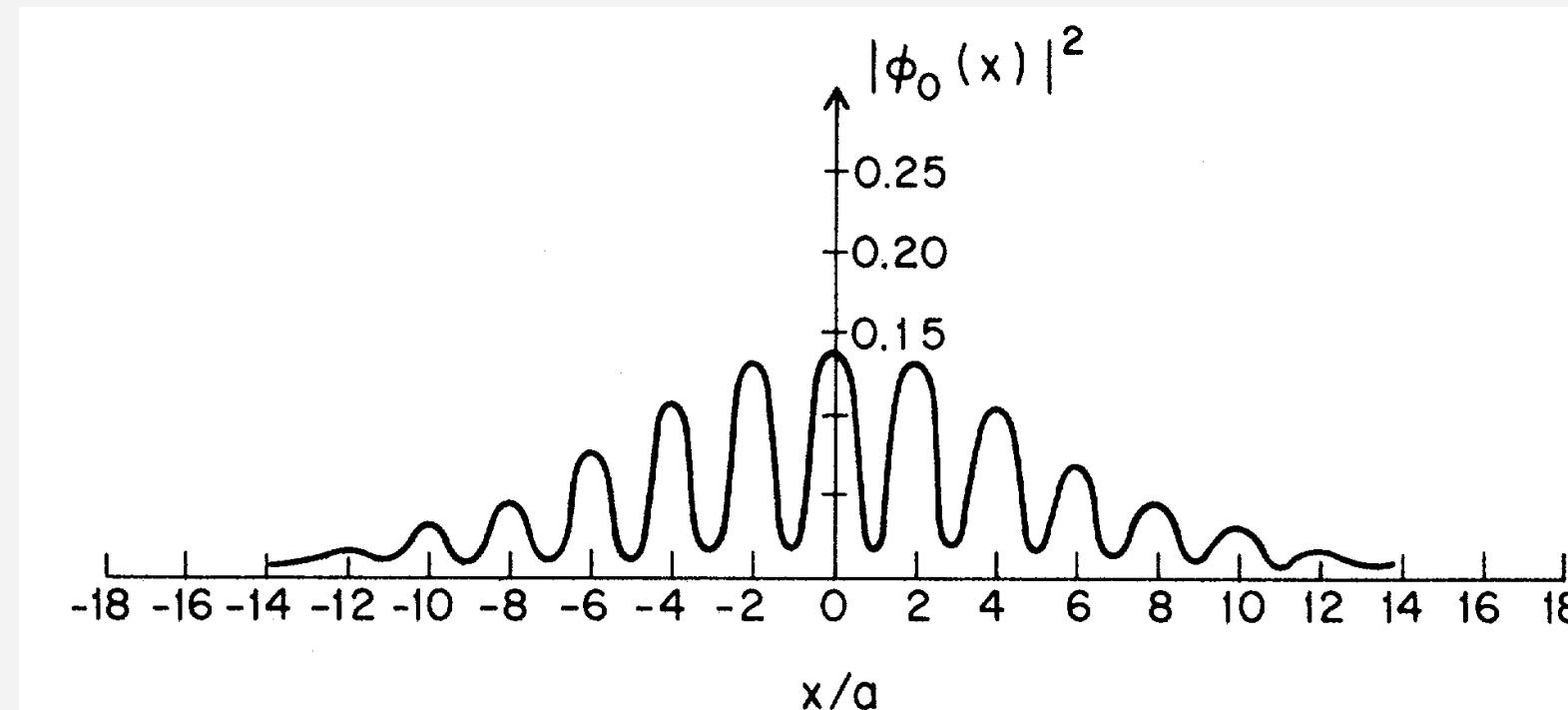
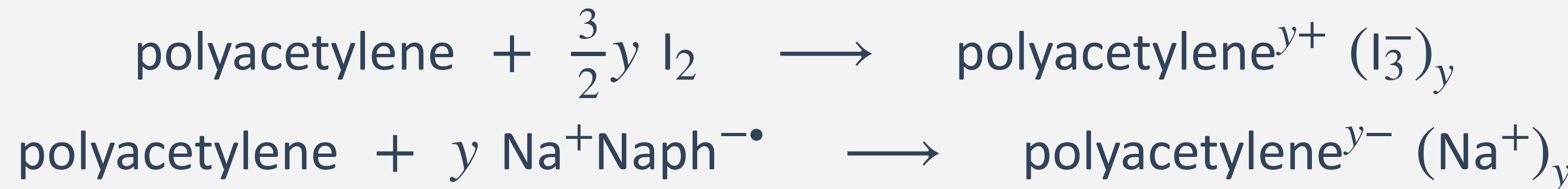


FIG. 3. Probability distribution of the localized electronic state at the center of the gap.

- Electron-Nuclear Double Resonance (ENDOR): soliton delocalized over 7–9 carbon atoms
- **Su-Schrieffer-Heeger (SSH) model: solitons are neutral defects with spin  $s = \frac{1}{2}$ , charge  $q = 0$**

# Chemical Doping

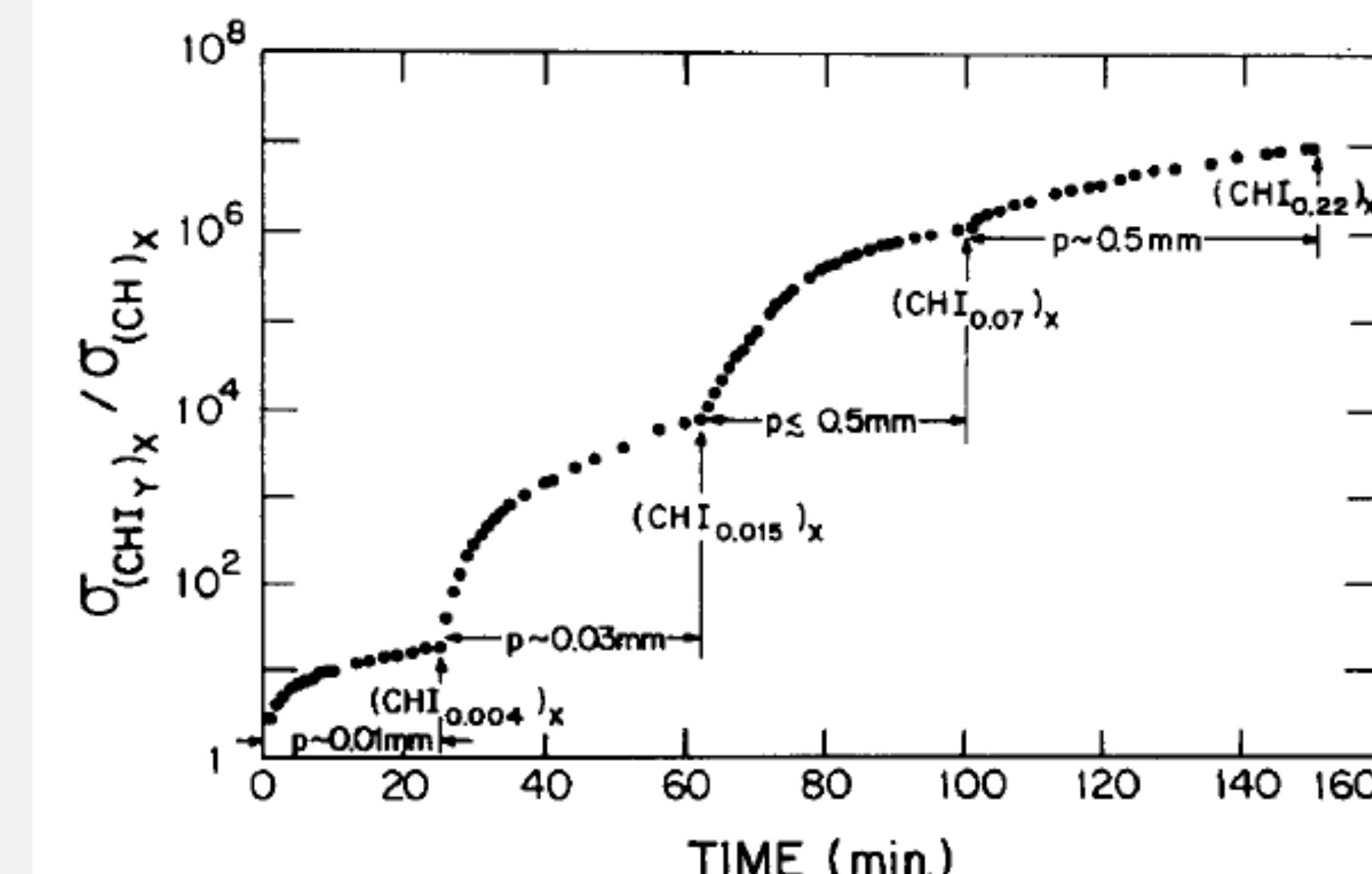
- doping of poly(acetylene) with halogens or alkali metals



**Table I.** Conductivity of Polycrystalline Polyacetylene and Derivatives (As-Grown Films)

Material	Conductivity, $\sigma (\Omega^{-1} \text{cm}^{-1})$ (25 °C)
<i>cis</i> -(CH) <sub>x</sub> <sup>a,b</sup>	$1.7 \times 10^{-9}$
<i>trans</i> -(CH) <sub>x</sub> <sup>a,b</sup>	$4.4 \times 10^{-5}$
<i>trans</i> -[(CH)(HBr) <sub>0.04</sub> ] <sub>x</sub>	$7 \times 10^{-4}$
<i>trans</i> -(CHCl <sub>0.02</sub> ) <sub>x</sub>	$1 \times 10^{-4}$
<i>trans</i> -(CHBr <sub>0.05</sub> ) <sub>x</sub> <sup>c</sup>	$5 \times 10^{-1}$
<i>trans</i> -(CHBr <sub>0.23</sub> ) <sub>x</sub> <sup>b,c</sup>	$4 \times 10^{-1}$
<i>cis</i> -[CH(ICI) <sub>0.14</sub> ] <sub>x</sub>	$5.0 \times 10^1$
<i>cis</i> -(CHI <sub>0.25</sub> ) <sub>x</sub>	$3.6 \times 10^2$
<i>trans</i> -(CHI <sub>0.22</sub> ) <sub>x</sub> <sup>b,c</sup>	$3.0 \times 10^1$
<i>trans</i> -(CHI <sub>0.20</sub> ) <sub>x</sub> <sup>b</sup>	$1.6 \times 10^2$
<i>cis</i> -[CH(IBr) <sub>0.15</sub> ] <sub>x</sub>	$4.0 \times 10^2$
<i>trans</i> -[CH(IBr) <sub>0.12</sub> ] <sub>x</sub>	$1.2 \times 10^2$
<i>trans</i> -[CH(AsF <sub>5</sub> ) <sub>0.03</sub> ] <sub>x</sub>	$7 \times 10^1$
<i>trans</i> -[CH(AsF <sub>5</sub> ) <sub>0.10</sub> ] <sub>x</sub> <sup>b</sup>	$4.0 \times 10^2$
<i>cis</i> -[CH(AsF <sub>5</sub> ) <sub>0.14</sub> ] <sub>x</sub>	$5.6 \times 10^2$
<i>trans</i> -[Na <sub>0.28</sub> (CH)] <sub>x</sub>	$8 \times 10^1$

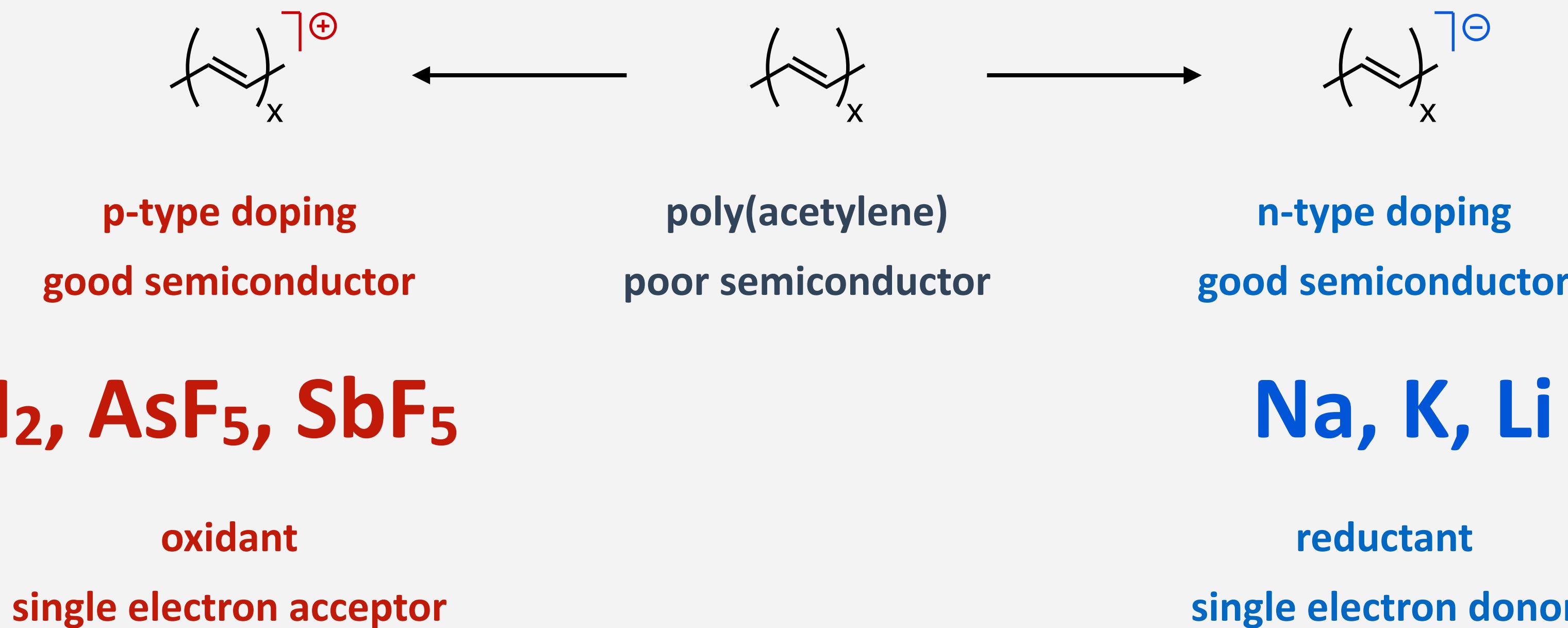
<sup>a</sup> H. Shirakawa, T. Ito, and S. Ikeda, unpublished results; see ref 1-3. <sup>b</sup> Composition obtained by chemical analysis from Galbraith Laboratories, Inc. (sum of all elements is  $\sim 99.8\text{--}100.1\%$ ). <sup>c</sup> See ref 1-3.



**FIG. 3.** Electrical conductivity (4-probe) of *trans*-(CH)<sub>x</sub> during doping reaction with iodine. The conductivity is normalized to that of the undoped sample.

- chemical doping of poly(acetylene) results in up to  $10^8$ -fold increase conductivity

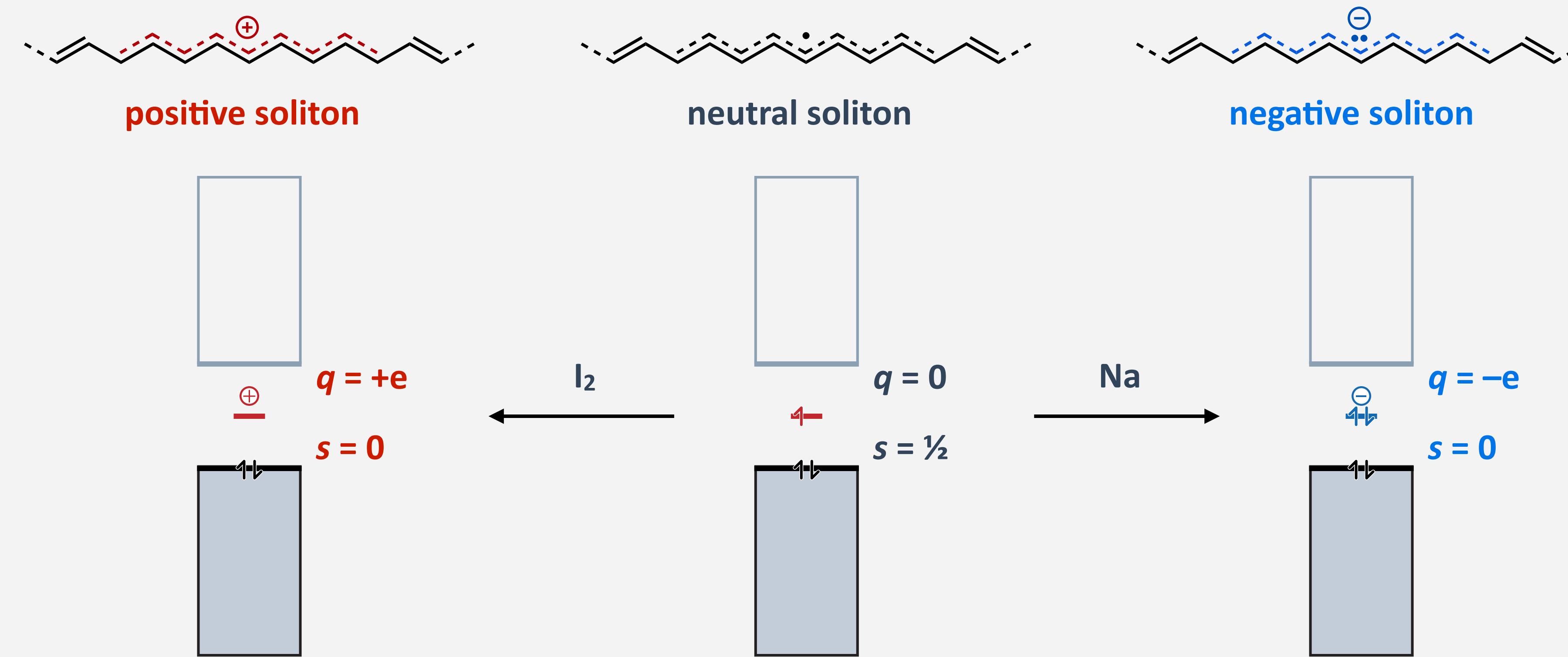
- chemically doped poly(acetylene) in the solid state becomes semiconducting or metallic



- different from inorganic semiconductor “impurity doping” (at ppm concentrations)
- “single electron transfer” oxidant/reductant, but must not induce follow-up reactions
- dopant applied at high concentrations (0.1–10 mol%), strongly disturbs structure/geometry
- conductivity & mobility increased by several orders of magnitude by chemical doping

# Formation of Positive or Negative Soliton Charge Carriers

- chemical doping converts neutral solitons into positive or negative soliton charge carriers

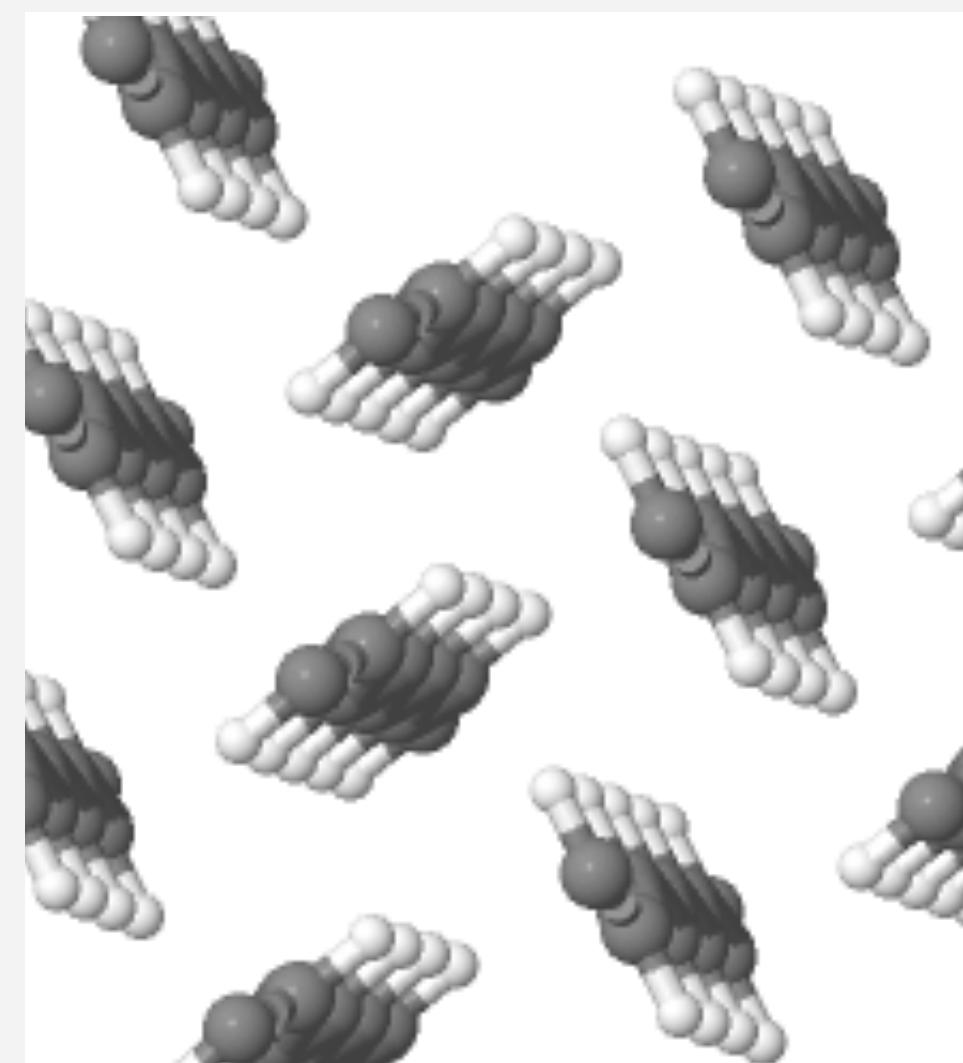


- positive/negative solitons have **no spin ( $s = 0$ )** but carry a **charge ( $q = \pm e$ )**
- organic chemistry view: carbocations or carbanions (delocalized over 7–23 carbon atoms)
- weakly doped poly(acetylene) is a semiconductor (isolated charges, limited delocalization)
- insulator-metal transition in strongly doped poly(acetylene) at about  **$10^{21}$  charges/cm $^{-1}$**

# Doping-Induced Phase Transition

- sub-stoichiometric amounts of dopant, unlike impurity doping in inorganic semiconductors

undoped polyacetylene



doping transition

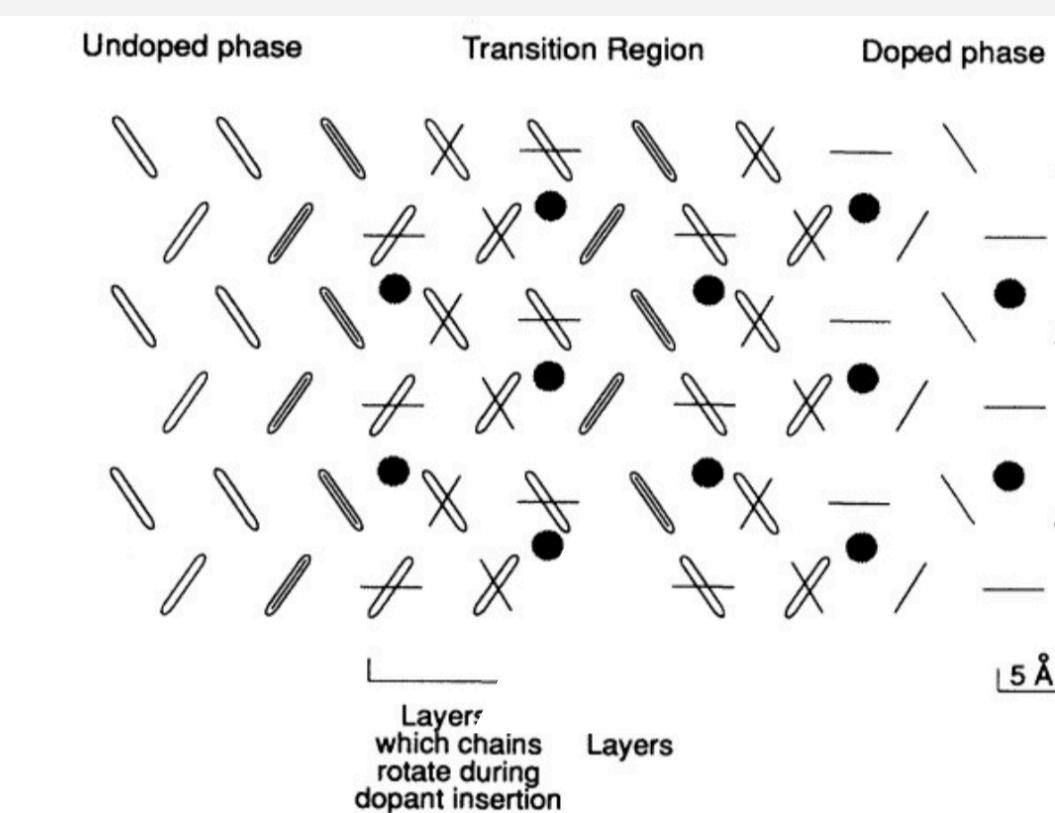


FIG. 4. Superposition of the lattices of undoped and lithium-doped polyacetylene in chain-axis projection. The open cigar-shaped boxes represent polyacetylene in the undoped phase. The lines represent the polyacetylene chains in the doped trigonal phase. The shaded circles are the  $\text{Li}^+$  ions.

Na-doped polyacetylene

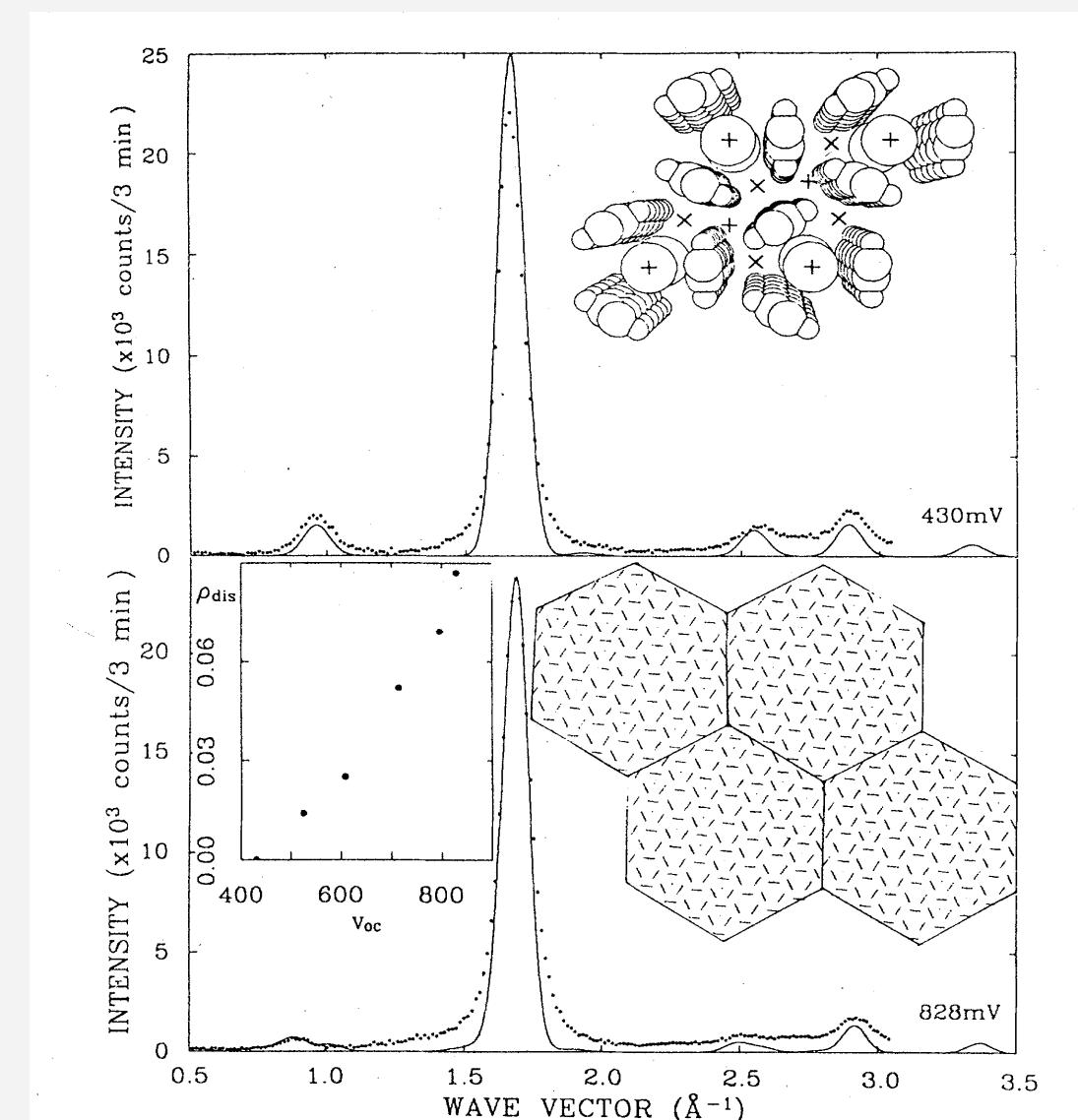


FIG. 7. Structure of Na-doped polyacetylene,  $[(\text{Na}^+)_y(\text{CH})^{+y}]_x$  (from Winokur *et al.*, 1987): (a) comparison between data (dots) and calculated diffraction profile (solid curve) for the commensurate  $\sqrt{3} \times \sqrt{3}$  superlattice channel structure,  $y \approx 0.11$ ; (b) comparison between data (dots) and calculated diffraction profile (solid curve) for discommensurate-domain structure shown in the inset,  $0.06 < y < 0.10$ .

- doping of polyacetylene results in phase transition to new materials with ion channels

# Charged Solitons

- charged solitons are supported by both theoretical models and experimental evidence

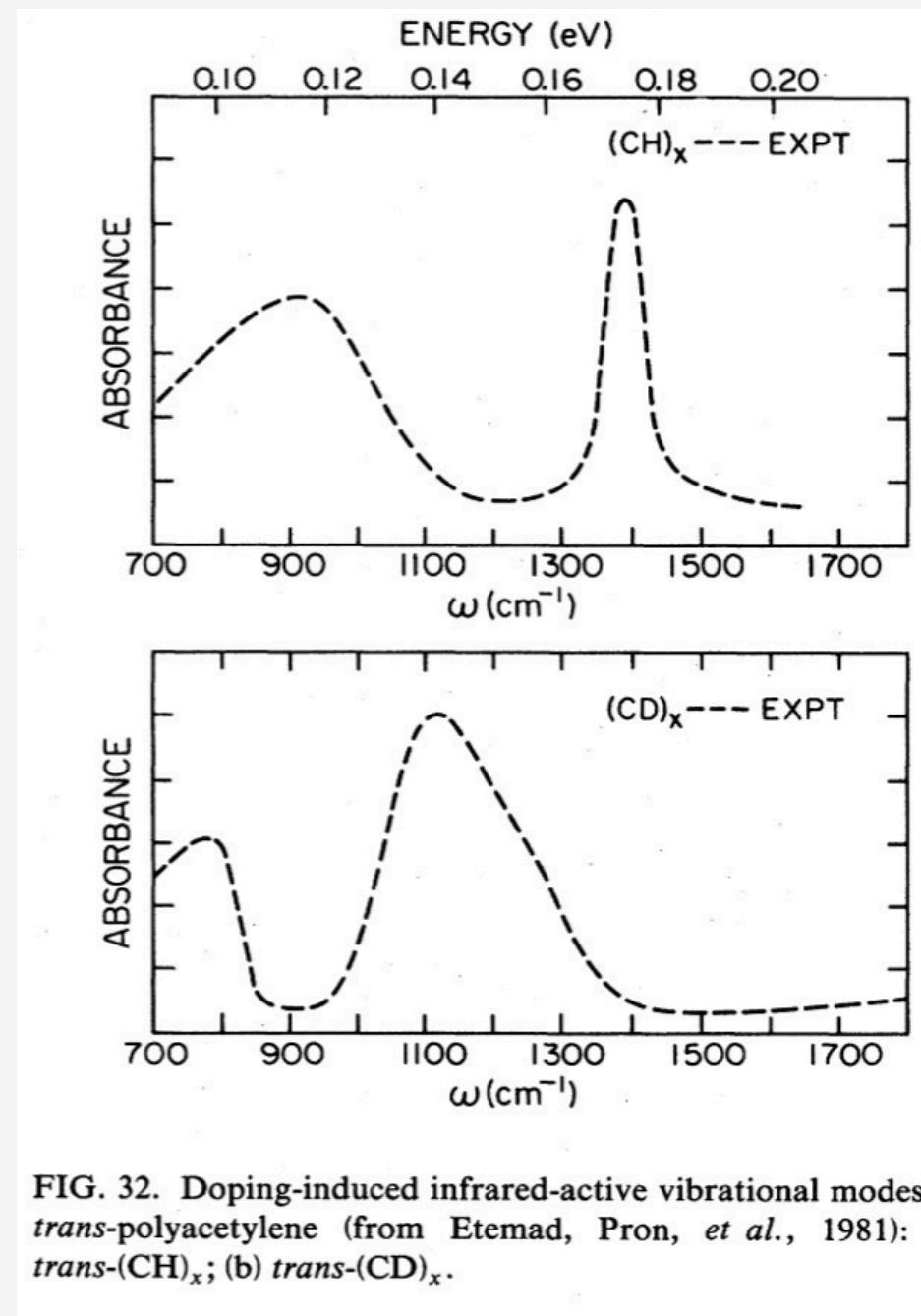


FIG. 32. Doping-induced infrared-active vibrational modes for *trans*-polyacetylene (from Etemad, Pron, *et al.*, 1981): (a) *trans*-(CH)<sub>x</sub>; (b) *trans*-(CD)<sub>x</sub>.

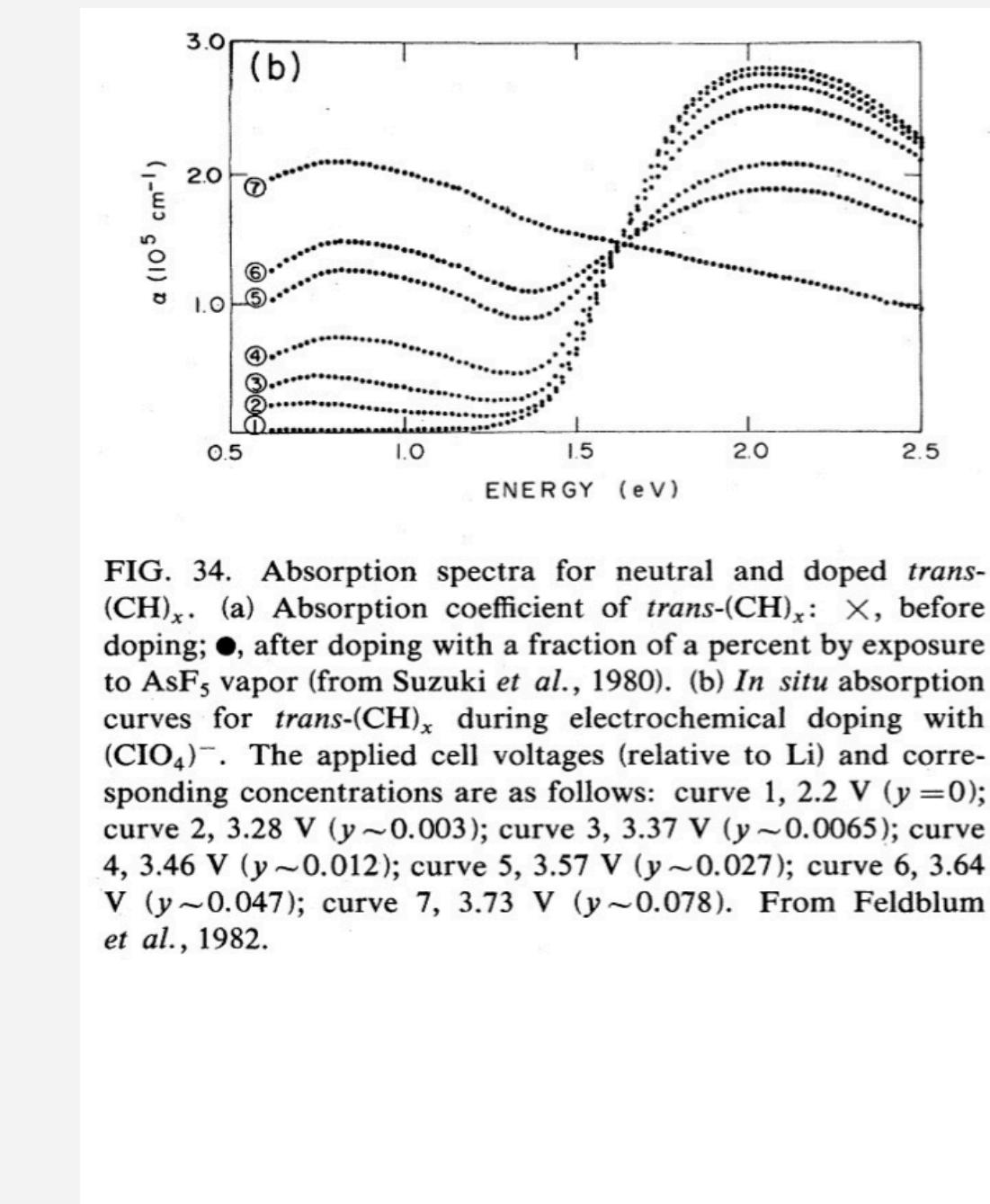


FIG. 34. Absorption spectra for neutral and doped *trans*-(CH)<sub>x</sub>. (a) Absorption coefficient of *trans*-(CH)<sub>x</sub>:  $\times$ , before doping;  $\bullet$ , after doping with a fraction of a percent by exposure to AsF<sub>5</sub> vapor (from Suzuki *et al.*, 1980). (b) *In situ* absorption curves for *trans*-(CH)<sub>x</sub> during electrochemical doping with (ClO<sub>4</sub>)<sup>-</sup>. The applied cell voltages (relative to Li) and corresponding concentrations are as follows: curve 1, 2.2 V ( $y=0$ ); curve 2, 3.28 V ( $y \sim 0.003$ ); curve 3, 3.37 V ( $y \sim 0.0065$ ); curve 4, 3.46 V ( $y \sim 0.012$ ); curve 5, 3.57 V ( $y \sim 0.027$ ); curve 6, 3.64 V ( $y \sim 0.047$ ); curve 7, 3.73 V ( $y \sim 0.078$ ). From Feldblum *et al.*, 1982.

IR: additional vibrational modes upon doping  
isotope-shift allows to link to modelling

UV-vis-NIR: appearance of absorptions within bandgap upon doping  
broad features implies mid-bandgap band

- charged soliton liquid undergoes first-order transition to a lattice of polaron-like defects